

Review of $\pi\pi$ scattering measurements in K decay

1. $\pi\pi$ scattering in $K \rightarrow 3\pi$ decay

- $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$

- Summary of results on $\pi\pi$ scattering lengths
- Comparison with the measurement of ponium lifetime

- $K_L \rightarrow \pi^0 \pi^0 \pi^0$

- Evidence for $\pi^+ \pi^- \rightarrow \pi^0 \pi^0$ in $K_L \rightarrow \pi^+ \pi^- \pi^0$ decay

- $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$

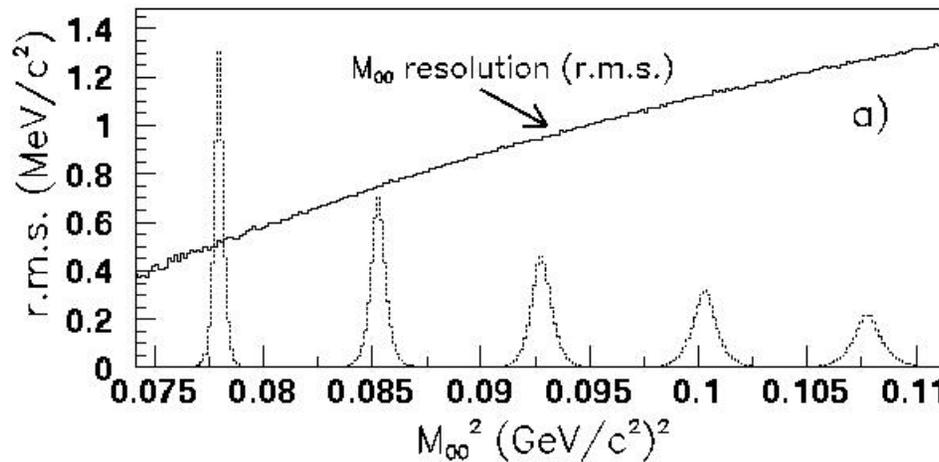
- Is $\pi\pi$ scattering observable by looking only at $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ decay ?

2. $\pi\pi$ scattering in $K^\pm \rightarrow \pi^+ \pi^- e^\pm (\bar{\nu})$ decays

3. Conclusions

$\pi\pi$ scattering in $K^\pm \rightarrow \pi^\pm\pi^0\pi^0$ decay

Observation of a cusp structure in the $\pi^0\pi^0$ invariant mass distribution at $M_{\pi^0\pi^0} = 2m_{\pi^+}$: an unexpected discovery from the NA48/2 high statistics study of $K^\pm \rightarrow \pi^\pm\pi^0\pi^0$ decay with good $M_{\pi^0\pi^0}$ resolution



$M_{\pi^0\pi^0}$ ($\equiv M_{00}$) resolution optimized for low M_{00} values

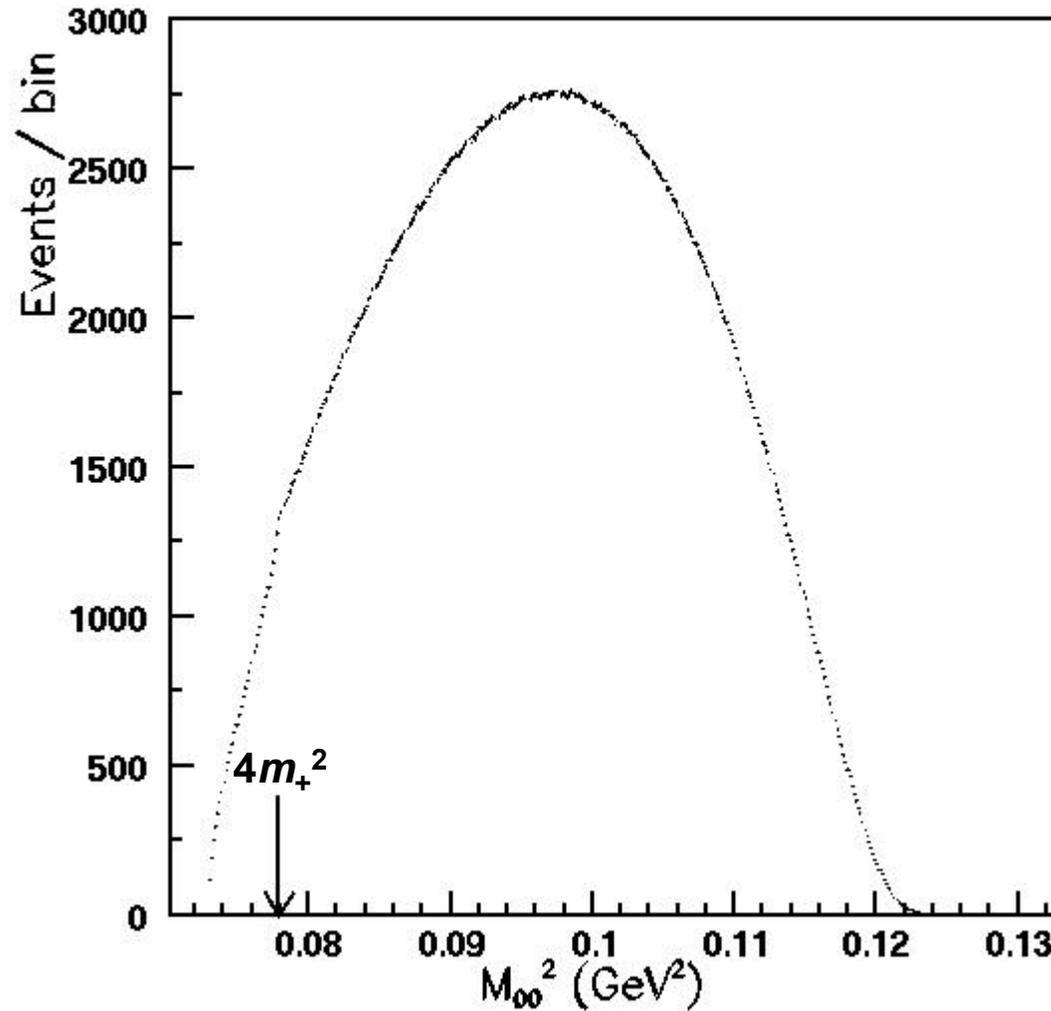
Initial motivation for this study:

the detection of $\pi^+\pi^-$ atom (“pionium”) formation in $K^\pm \rightarrow \pi^\pm\pi^+\pi^-$ decays, followed by pionium annihilation to $\pi^0\pi^0$

M_{00} resolution $\sigma = 0.56 \text{ MeV}$ at $M_{00} = 4m_+$

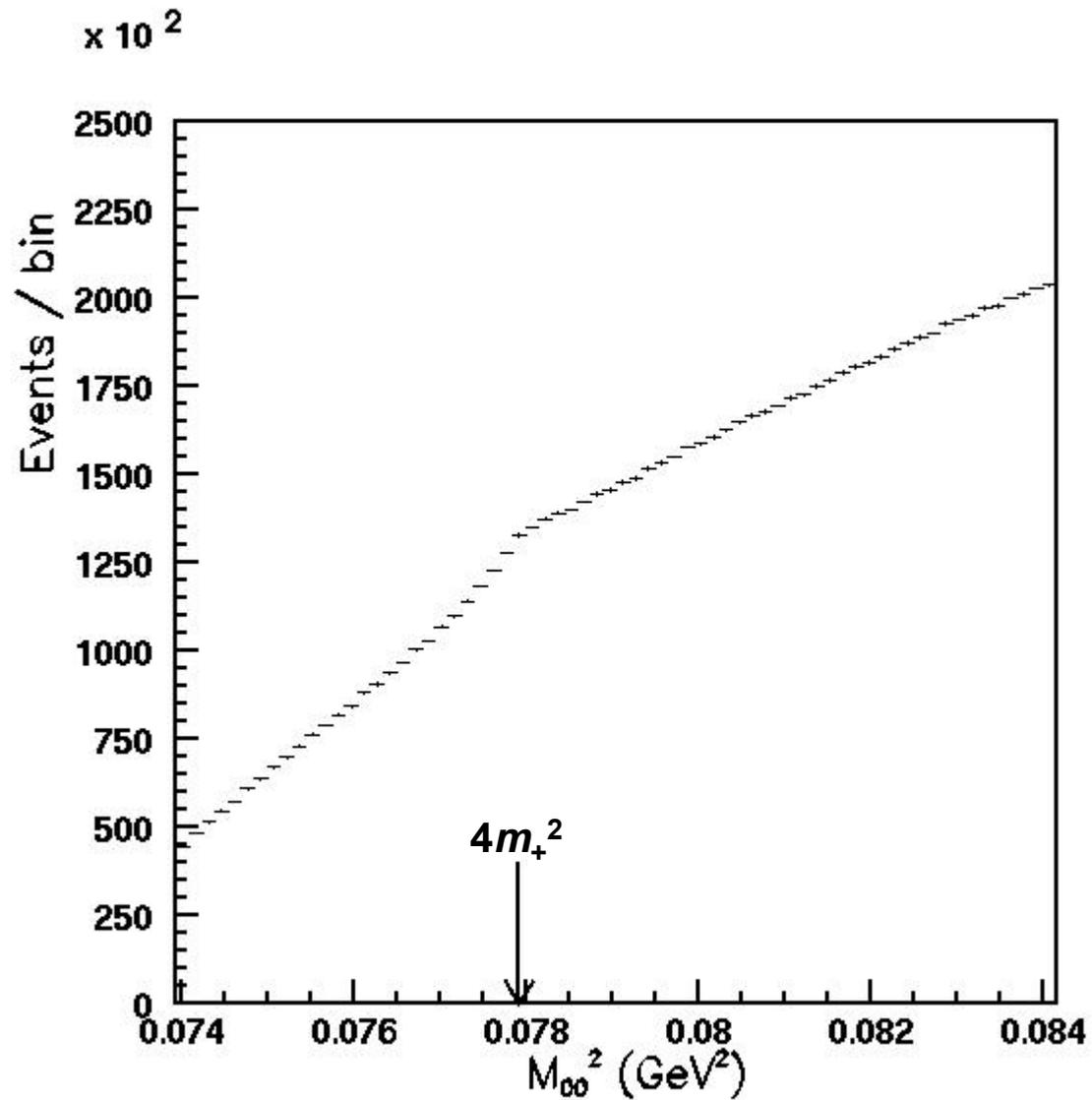
NA48/2 (PRELIMINARY)

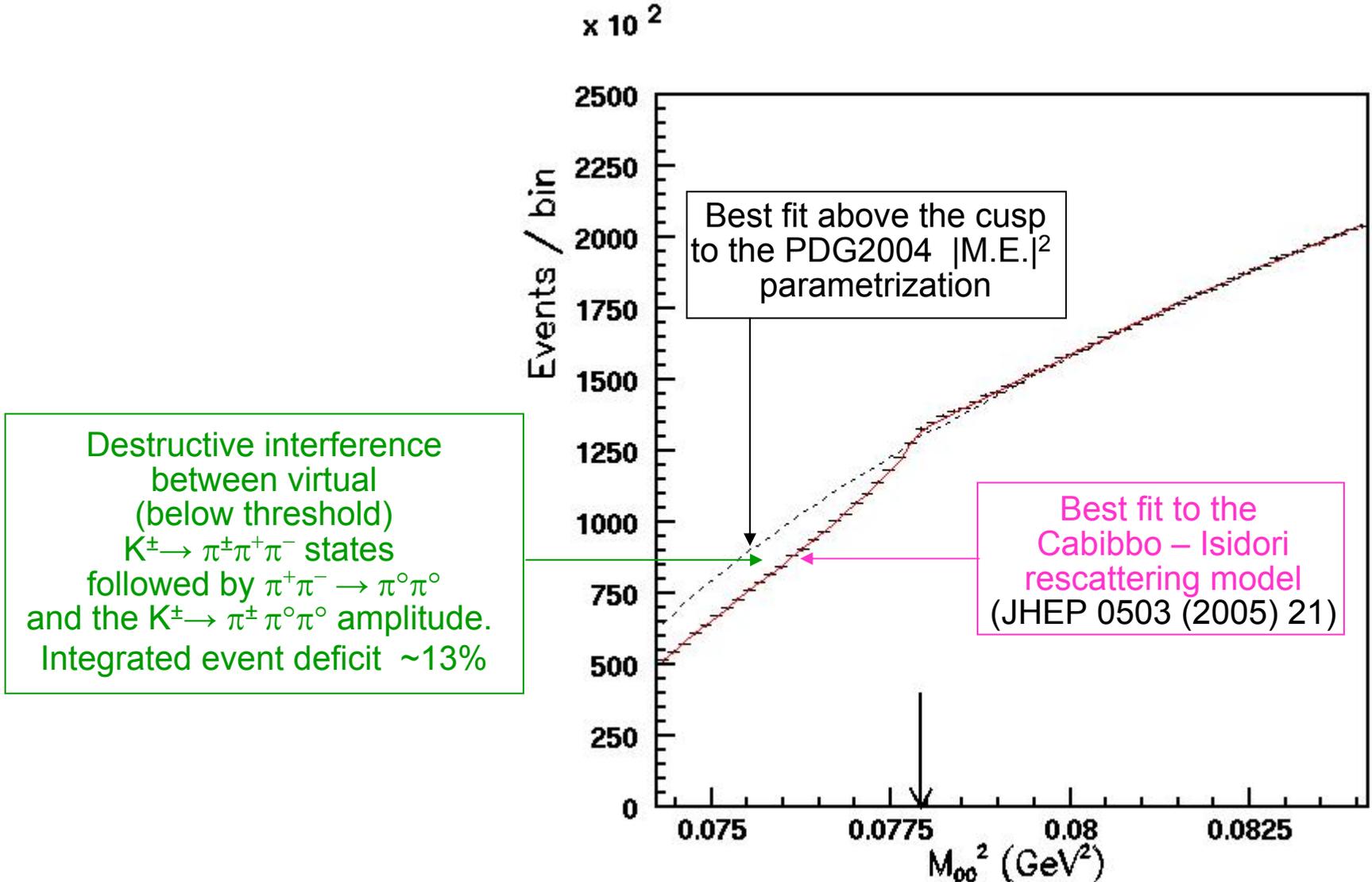
59,624,170 fully reconstructed $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ events
 $\times 10^2$



$$m_+ \equiv m(\pi^+)$$

ZOOM ON THE CUSP REGION





The $\pi^+ \pi^-$ charge exchange amplitude near threshold is proportional to the difference of $\pi\pi$ scattering lengths $a_0 - a_2$

➔ a new method to measure $a_0 - a_2$ and its sign

N. Cabibbo, PRL 93 (2004) 12181

Fit results:

$$(a_0 - a_2)m_+ = 0.261 \pm 0.006 \pm 0.003 \pm 0.0013 \pm 0.013$$

(stat.) (syst.) (ext.) (theor.)

$$a_2 m_+ = -0.037 \pm 0.013 \pm 0.009 \pm 0.002$$

(the sensitivity to a_2 comes from higher-order terms)

External uncertainty:

from the uncertainty on the ratio of $K^+ \rightarrow \pi^+ \pi^+ \pi^-$ and $K^+ \rightarrow \pi^+ \pi^0 \pi^0$ decay widths:

$$\frac{\Gamma(K^+ \rightarrow \pi^+ \pi^+ \pi^-)}{\Gamma(K^+ \rightarrow \pi^+ \pi^0 \pi^0)} = 3.182 \pm 0.047 \quad \text{(PDG 2006)}$$

giving $\frac{A(K^+ \rightarrow \pi^+ \pi^+ \pi^-)}{A(K^+ \rightarrow \pi^+ \pi^0 \pi^0)} = 1.975 \pm 0.015$ at the Dalitz plot centres ($u = v = 0$)

(exact isospin symmetry predicts 2)

Theoretical uncertainty on $(a_0 - a_2)m_+$: $\pm 5\%$

(estimated effect from neglecting higher order diagrams and radiative corrections)

Fit with analyticity and chiral symmetry constraint

between a_0 and a_2 (Colangelo, Gasser, Leutwyler, PRL 86 (2001) 5008)

$$(a_0 - a_2)m_+ = 0.263 \pm 0.003 \pm 0.0014 \pm 0.0013 \pm 0.013$$

(stat.) (syst.) (ext.) (theor.)

GOOD AGREEMENT WITH THE RESULT OF THE DIRAC EXPERIMENT :

Pionium mean lifetime $\tau_{1s} = (2.91^{+0.24}_{-0.43}) \times 10^{-15} \text{ s}$

$\Rightarrow |a_0 - a_2| m_+ = 0.264^{+0.020}_{-0.011}$

**Central values
= old values !**

NA48/2 : $(a_0 - a_2)m_+ = 0.261 \pm 0.006 \pm 0.003 \pm 0.0013 \pm 0.013$
stat. syst. ext. theor.

Very little theoretical uncertainty in the prediction of the pionium lifetime because the interaction responsible for $\pi^+\pi^- \rightarrow \pi^0\pi^0$ is made effectively “weak” by the large pionium radius:

$$R_{\text{pionium}} \approx R_\infty \frac{2m_e}{m_+} \approx 3.9 \times 10^{-11} \text{ cm} \quad (R_\infty : \text{Bohr radius for } M_{\text{nucleus}} = \infty)$$

$R_{\text{pionium}} \gg$ strong interaction radius ($\sim 10^{-13} \text{ cm}$)

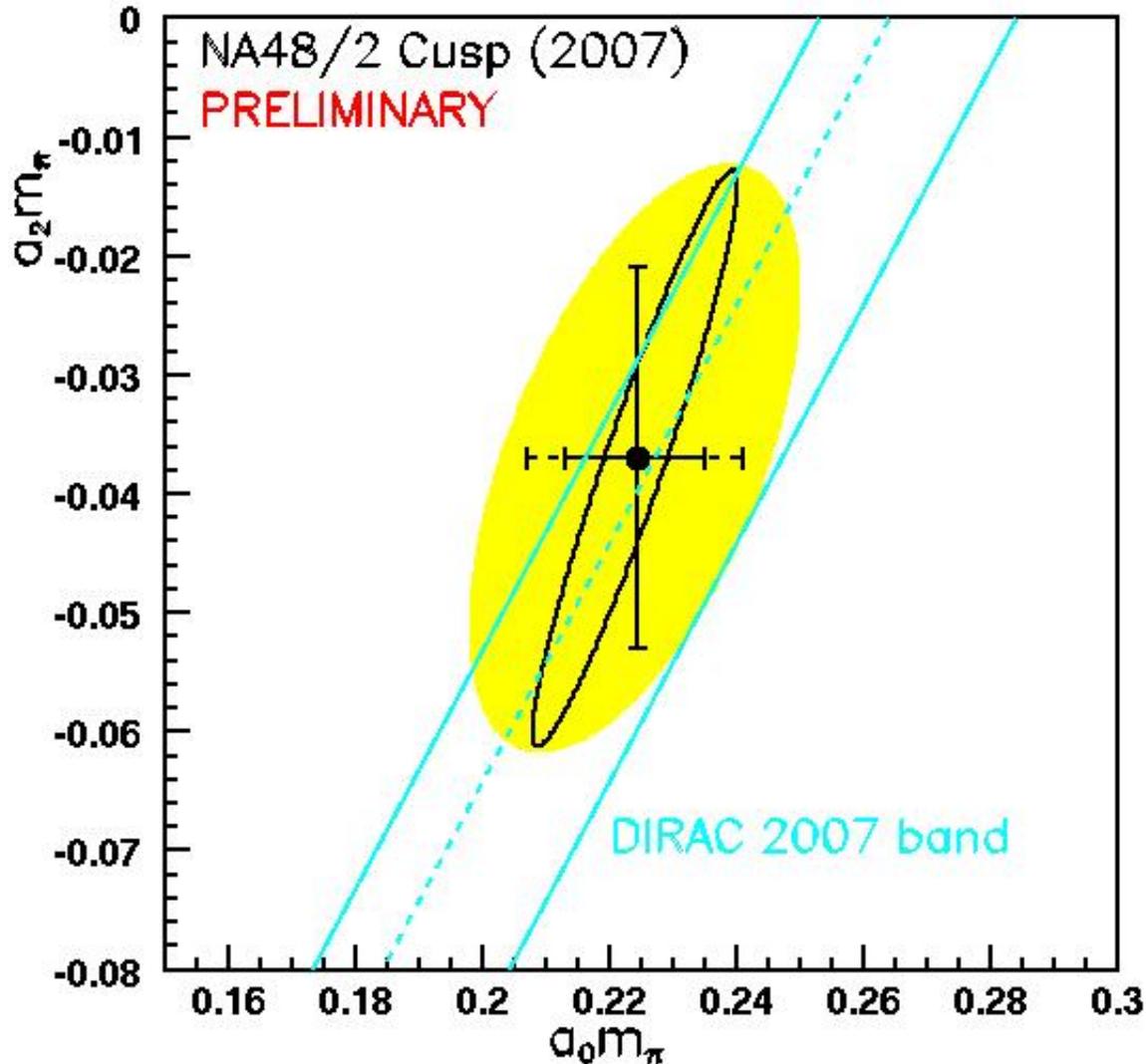
\Rightarrow very little overlap of the $\pi^+\pi^-$ atomic wave function with the strong interaction volume

NA48/2 (PRELIMINARY): from (a0 – a2) and a2 extract a0

(must take into account the statistical error correlation coefficient ≈ -0.92)

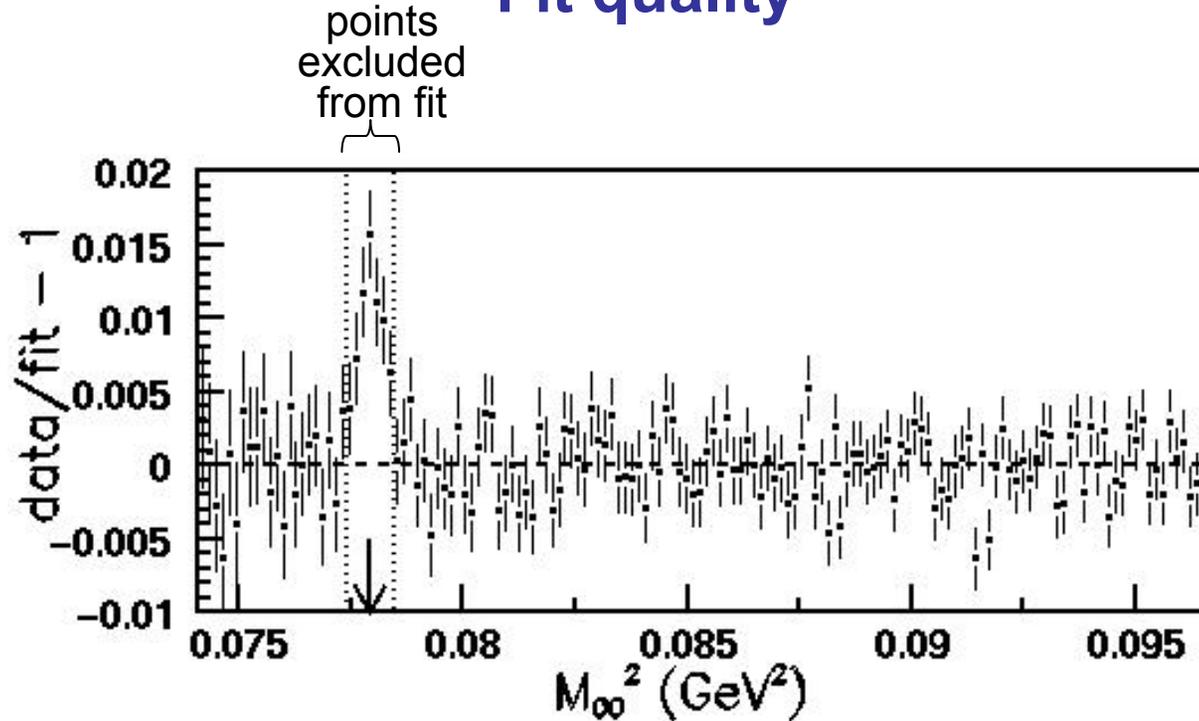
$$a_0 m_\pi = 0.224 \pm 0.008 \pm 0.006 \pm 0.003 \pm 0.013$$

stat. svst. ext. theor.



The yellow area represents theoretical uncertainty (assumed Gaussian)
The dashed bars represent the theoretical uncertainty

Fit quality



Observe an excess of events in the M_{00}^2 interval excluded from fit.
Formation of $\pi^+\pi^-$ atoms (pionium) decaying to $\pi^0\pi^0$?

Rate = $(1.82 \pm 0.21) \times 10^{-5}$ per $K^\pm \rightarrow \pi^\pm\pi^+\pi^-$ decay

**= 2.28 ± 0.26 times the predicted pionium formation rate
(= 0.8×10^{-5} , Silagadze 1994)**

(adding a δ – function at $M_{00} = 2m_+$ to the theoretical spectrum before taking into account detector effects, and using its integral as a free parameter in the fit)

Direct count of the event excess in the bins excluded from the fit:

Rate = 2.58 ± 0.31 times the predicted pionium formation rate

(960 ± 320 events more than the pionium best fit)

IS THE EVENT EXCESS ONLY DUE TO PIONIUM ATOMS?

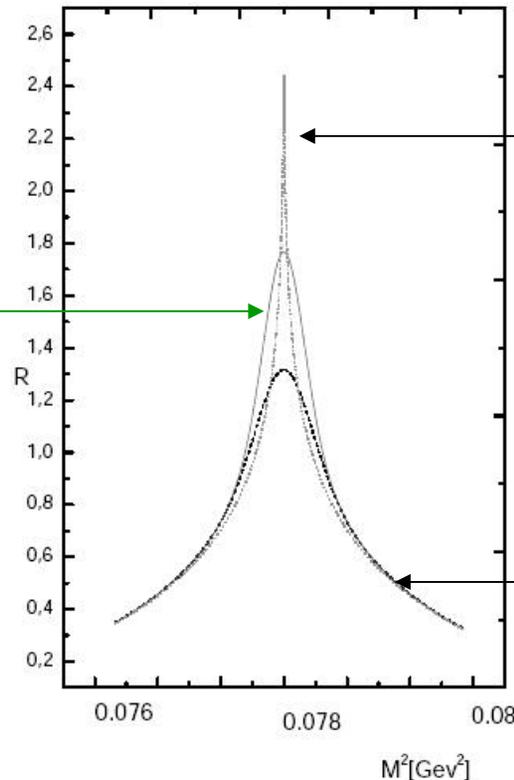
Electromagnetic corrections to final state interactions in $K \rightarrow 3\pi$ decays

(Gevorkian, Tarasov, Voskresenskaya, hep-ph / 0612129)

Two contributions from $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ decay to the $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ cusp region:

- Pionium formation : $\pi^+ \pi^-$ atom $\rightarrow \pi^0 \pi^0$ (negligible width)
- Additional $\pi^+ \pi^-$ unbound states with resonance structure $\rightarrow \pi^0 \pi^0$

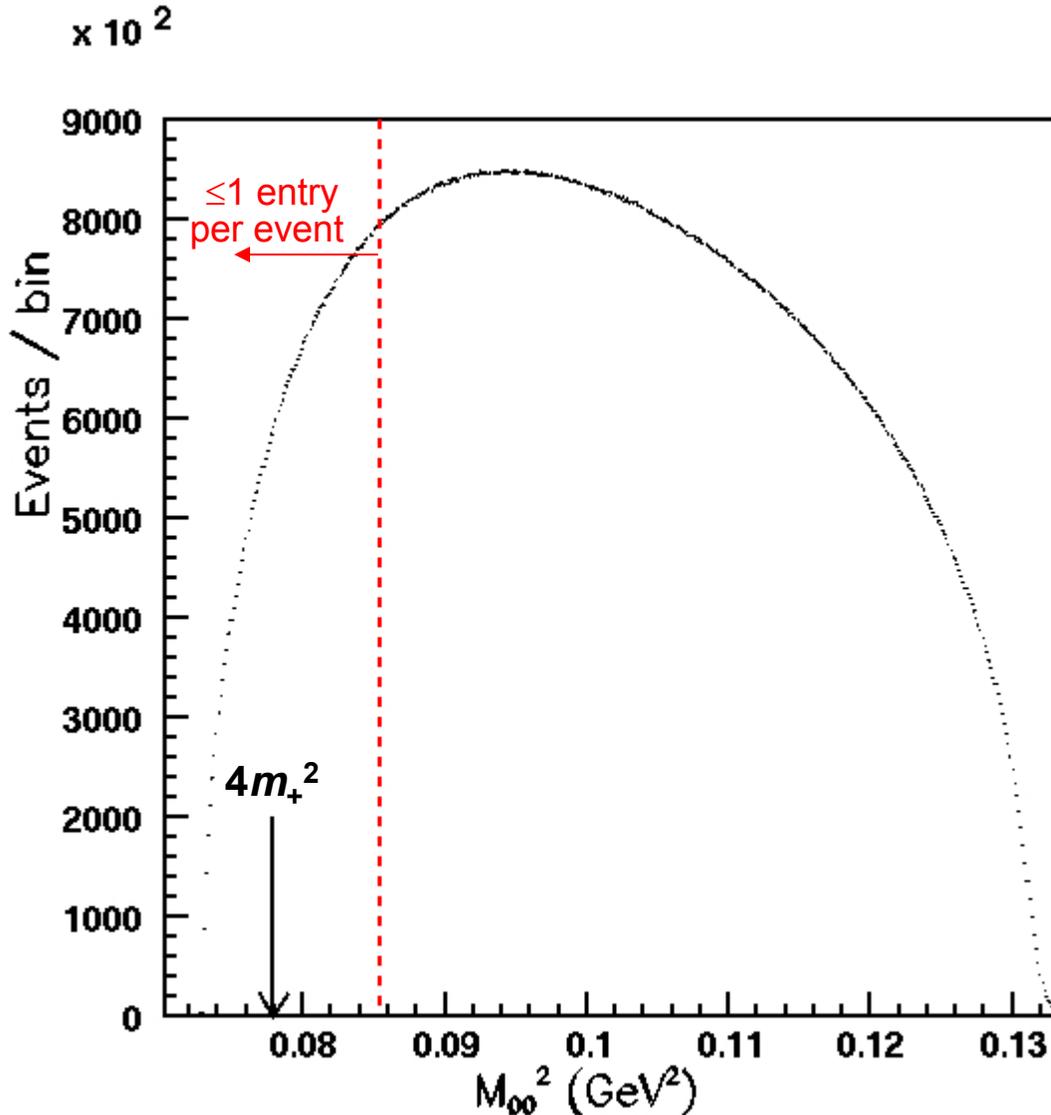
$\pi^+ \pi^-$ atoms and
 $\pi^+ \pi^-$ resonant structure
with experimental resolution



$\pi^+ \pi^-$ resonant structure
(no experimental resolution)

$\pi^+ \pi^-$ resonant structure
with experimental resolution

$K_L \rightarrow \pi^0 \pi^0 \pi^0$
87,950,601 events (NA48)
(3 entries per event)



**CUSP STRUCTURE
MUCH LESS VISIBLE
THAN IN $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ DECAY**

$K^+ \rightarrow \pi^+\pi^0\pi^0$, $K_L \rightarrow \pi^0\pi^0\pi^0$ comparison

$$\mathcal{R}(K^+) \approx \frac{2\mathcal{M}_{++-}\mathcal{M}_{+00}}{(\mathcal{M}_{+00})^2} = \frac{2\mathcal{M}_{++-}}{\mathcal{M}_{+00}} \quad \left\{ \begin{array}{l} \mathcal{M}_{++-} : K^+ \rightarrow \pi^+\pi^+\pi^- \text{ matrix element} \\ \mathcal{M}_{+00} : K^+ \rightarrow \pi^+\pi^0\pi^0 \text{ matrix element} \end{array} \right.$$

two possible $\pi^+\pi^-$ pairs

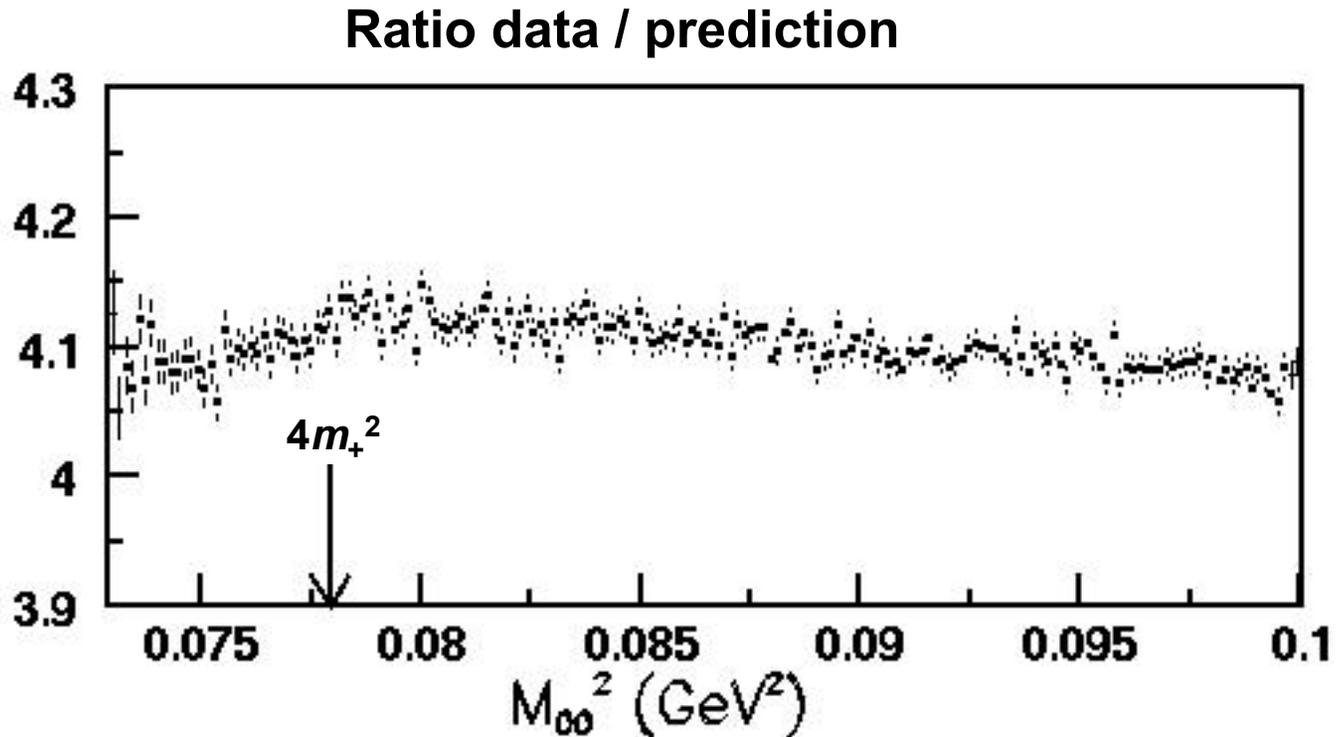
$$\mathcal{R}(K_L) \approx \frac{\mathcal{M}_{+-0}\mathcal{M}_{000}}{(\mathcal{M}_{000})^2} = \frac{\mathcal{M}_{+-0}}{\mathcal{M}_{000}} \quad \left\{ \begin{array}{l} \mathcal{M}_{+-0} : K_L \rightarrow \pi^+\pi^-\pi^0 \text{ matrix element} \\ \mathcal{M}_{000} : K_L \rightarrow \pi^0\pi^0\pi^0 \text{ matrix element} \end{array} \right.$$

Calculate matrix elements at cusp point ($M_{\pi\pi} = 2m_+$)
from measured partial width ratios and slope parameters:

$$\mathcal{R}(K^+) \approx 6.1 ; \mathcal{R}(K_L) \approx 0.47 \quad \longrightarrow \quad \frac{\mathcal{R}(K^+)}{\mathcal{R}(K_L)} \approx 13$$

Cusp “visibility” is ~ 13 times higher in $K^+ \rightarrow \pi^+\pi^0\pi^0$ decays
than in $K_L \rightarrow \pi^0\pi^0\pi^0$ decays

Compare data with predicted distribution without $\pi^+\pi^-$ charge exchange
(obtained from Monte Carlo simulation)



➔ Evidence for a change of slope in the data
near the expected cusp point

Rescattering in $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ decay

First attempt to perform a combined fit of $M(\pi\pi)$ distributions from $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ and $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ decays

using the Bern – Bonn non-relativistic effective field theory

(Colangelo, Gasser, Kubis, Rusetsky, Phys. Lett. B 638 (2006) 187)

(VERY PRELIMINARY – only statistical errors are shown)

Fit with 8 free parameters:

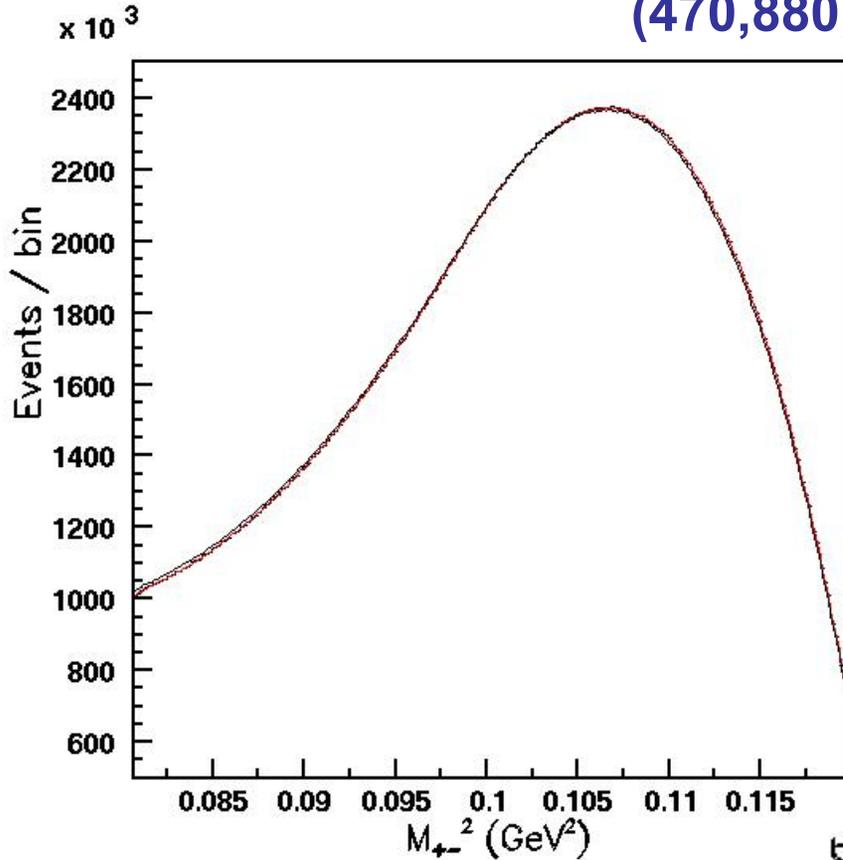
- 1 normalization constant, 2 slope parameters (g_0, h_0) for $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ M_{00}^2 distribution (2004 data) ; $k_0 = 0.0096$ is kept fixed ;
- 1 normalization constant common to the M_{+-}^2 and $M_{\pm\pm}^2$ distributions, 3 slope parameters (g, h, k) for $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ (2003 data);
- $(a_0 - a_2)m_+$; $a_2 m_+ = -0.044$ (fixed in the fit)

Fit intervals:

- $0.074094 < M_{00}^2 < 0.11744 \text{ GeV}^2$
- $0.080694 < M_{+-}^2, M_{\pm\pm}^2 < 0.119844 \text{ GeV}^2$

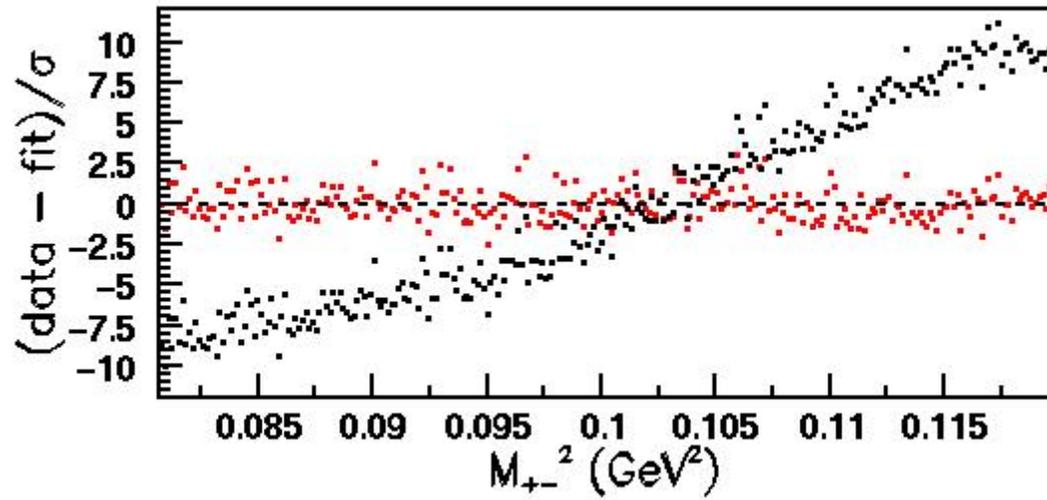
$$\text{Fit results} \left\{ \begin{array}{l} \chi^2 = 757.1 / 757 \text{ d.f.} \\ g_0 = 0.6191 \pm 0.0008 ; h_0 = -0.0497 \pm 0.0006 \\ g = -0.1837 \pm 0.0006 ; h = 0.0004 \pm 0.0003 ; k = -0.0059 \pm 0.0001 \\ (a_0 - a_2)m_+ = 0.266 \pm 0.003 \end{array} \right.$$

M_{+-}^2 distribution from $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ decay (470,880,352 entries)

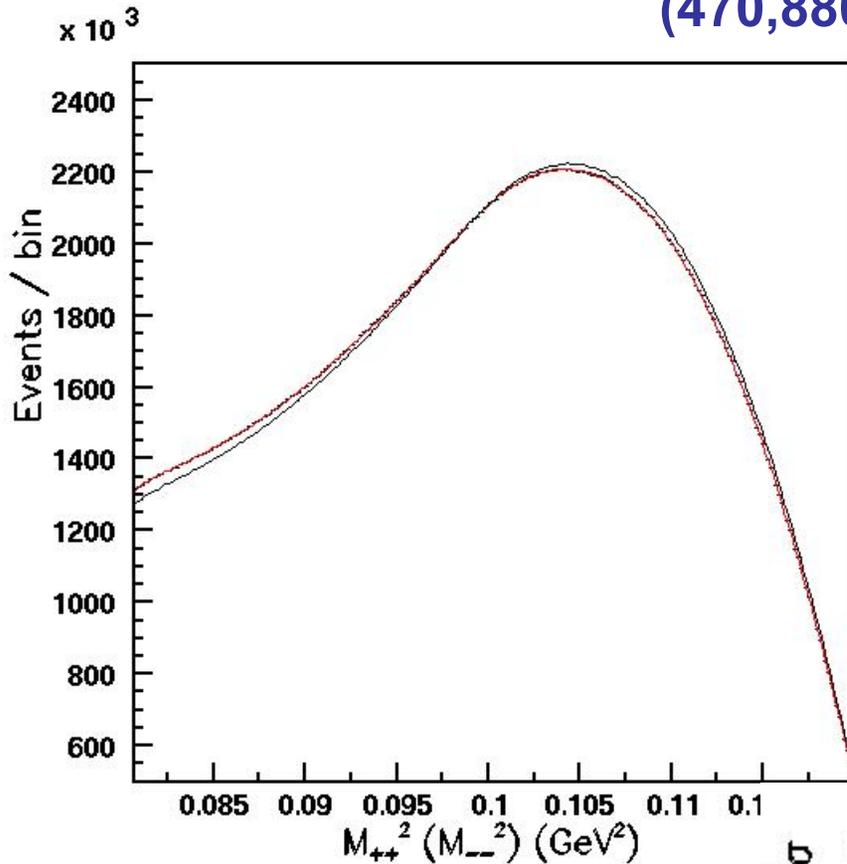


— Combined best fit
with rescattering

— Best fit to $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ only;
no rescattering;
slope parameter values
as from fit with rescattering.
 $\chi^2 = 9573 / 260$ d.f.

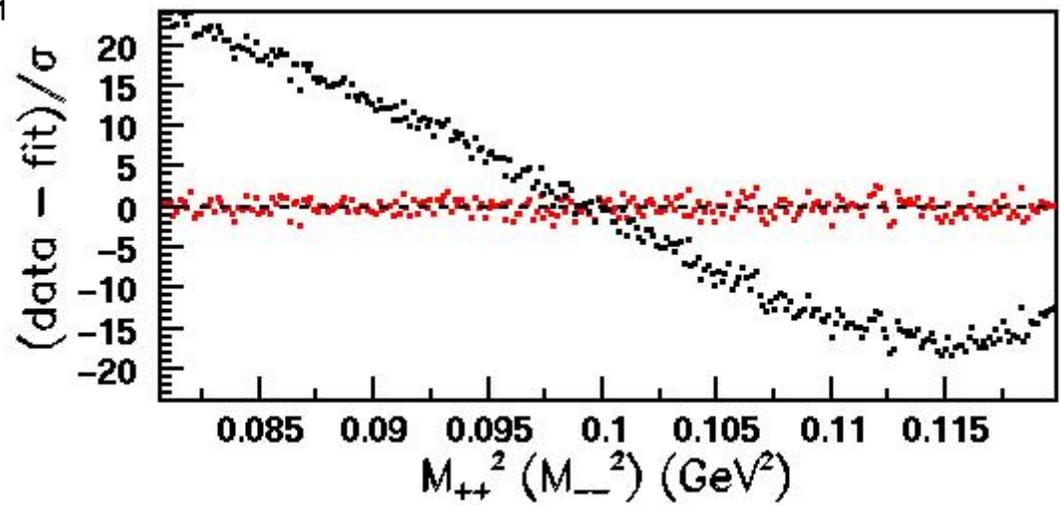


$M_{\pm\pm}^2$ distribution from $K^\pm \rightarrow \pi^\pm\pi^+\pi^-$ decay (470,880,352 events)

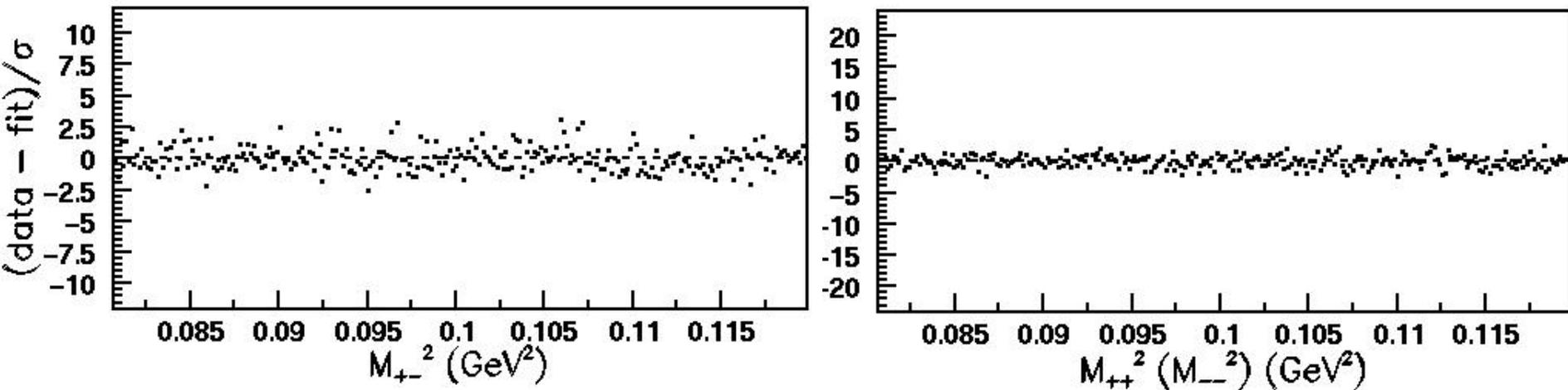


— Combined best fit
with rescattering

— Best fit to $K^\pm \rightarrow \pi^\pm\pi^+\pi^-$ only;
no rescattering;
slope parameter values
as from fit with rescattering.
 $\chi^2 = 47106 / 260$ d.f.



Fit to $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ only
No rescattering;
slope parameters allowed to vary
 $\chi^2 = 516.9 / 517$ d.f.



- Rescattering effects in $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ decay can be “faked” by different slope parameter values:

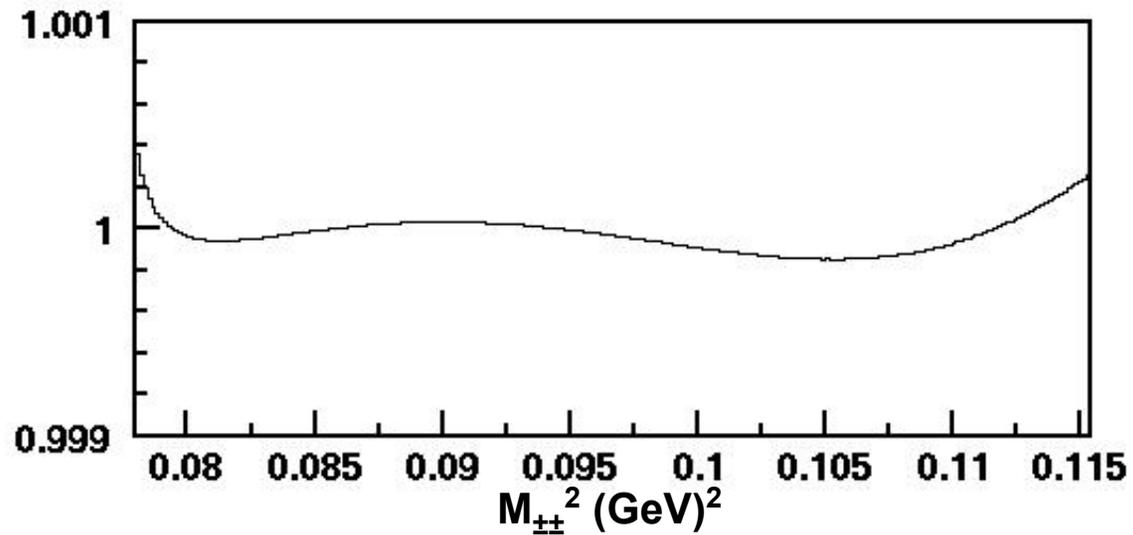
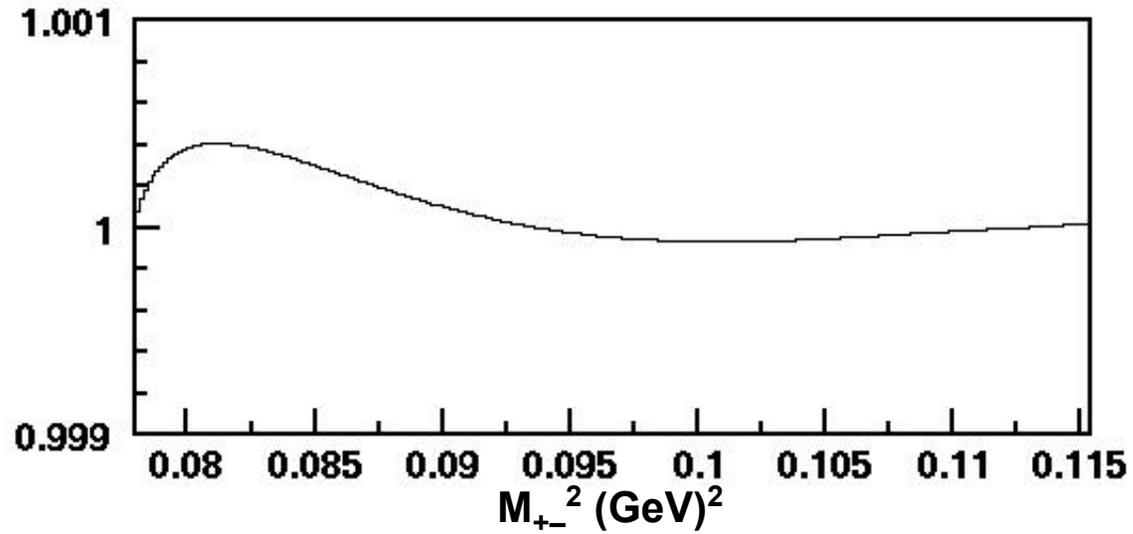
$$g = -0.2101 \pm 0.0001$$

$$h = 0.0088 \pm 0.0003$$

$$k = -0.0039 \pm 0.0001$$

- It is very difficult (impossible?) to obtain evidence for $\pi\pi$ rescattering from $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ decay alone

Ratio $\frac{\text{Best fit with rescattering}}{\text{Best fit without rescattering}}$



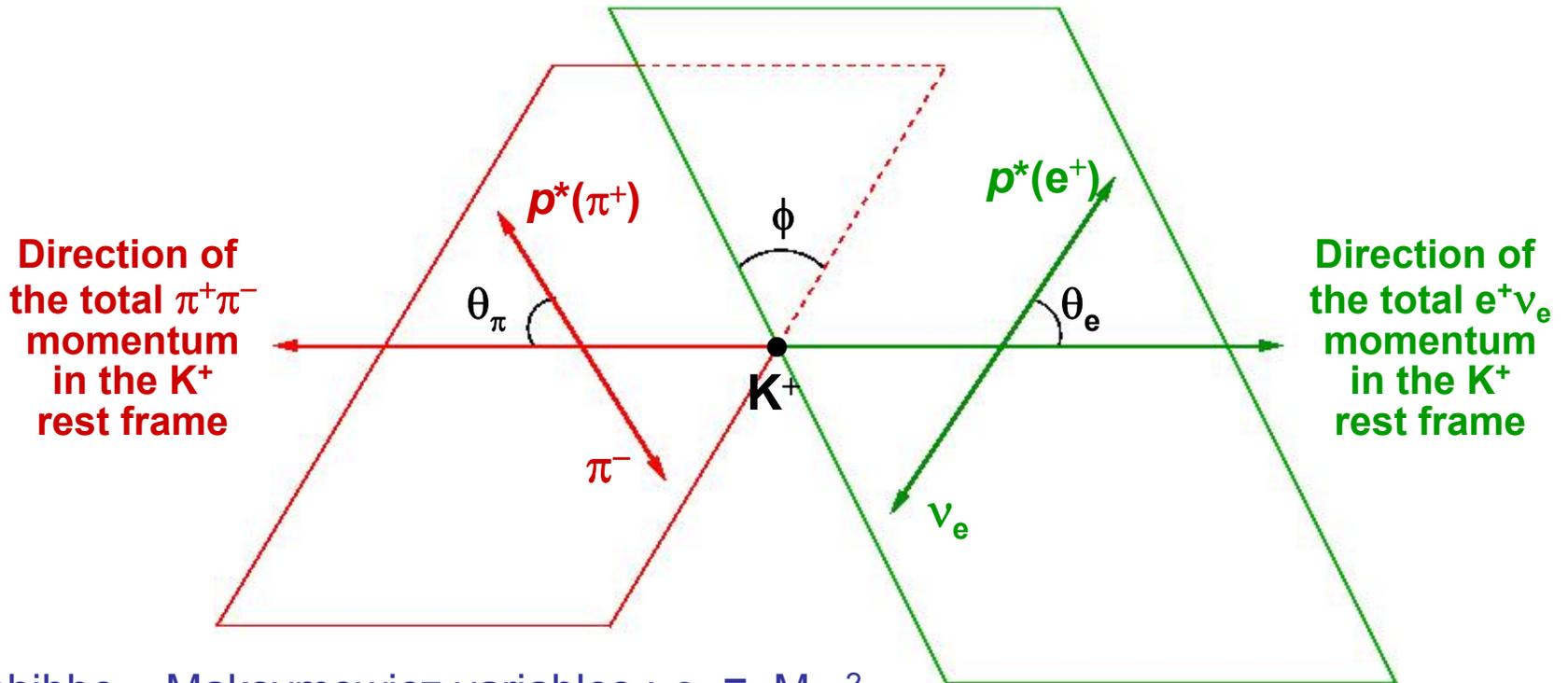
A message to the Particle Data Group

The parametrization of the $|\text{matrix element}|^2$ for $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$, $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$, $K_L \rightarrow \pi^+ \pi^- \pi^0$, and $K_L \rightarrow \pi^0 \pi^0 \pi^0$ by a series expansion of two Dalitz variables (see PDG 2006, 660-1, 681) should be abandoned:

- It cannot be used to include other amplitudes interfering with the main, weak decay amplitude;
- It does not describe the cusp structure observed in $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ by NA48/2;
- The slope parameter values (coefficients of the series expansions) depend on whether $\pi\pi$ scattering in the final state is taken into account.

$\pi\pi$ scattering in $K^\pm \rightarrow \pi^+\pi^-e^\pm \bar{\nu}$ decay

A rare decay [B.R. = $(4.09 \pm 0.09) \times 10^{-5}$] described by five independent variables



Cabibbo – Maksymowicz variables :

$$s_\pi \equiv M_{\pi\pi}^2$$

$$s_e \equiv M_{e\nu}^2$$

$$\theta_e$$

$$\theta_\pi$$

$$\phi$$

For $K^+ \Rightarrow K^-$
 $\phi \Rightarrow \pi + \phi$
 $\theta_e \Rightarrow \pi - \theta_e$

The Ke4 decay amplitude depends on two complex phases identified with

- δ_0 : the $\pi\pi$ scattering phase shift in the $I = 0$, $\ell = 0$ state (s – wave)
- δ_1 : the $\pi\pi$ scattering phase shift in the $I = 1$, $\ell = 1$ state (p – wave)

The Ke4 decay rate depends on the phase shift difference $\delta = \delta_0 - \delta_1$

δ is an increasing function of $M_{\pi\pi}$; $\delta \rightarrow 0$ for $M_{\pi\pi} \rightarrow 2m_+$

(from scattering theory : $\delta(k) \approx ak$ at very low centre-of-mass momentum k ;
 a is the scattering length ; $a \neq 0$ for s – waves only)

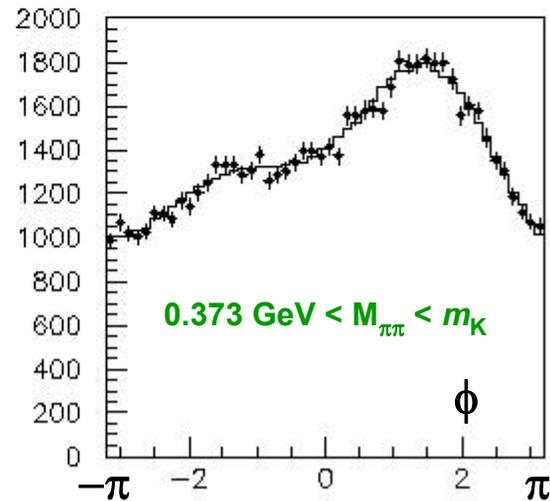
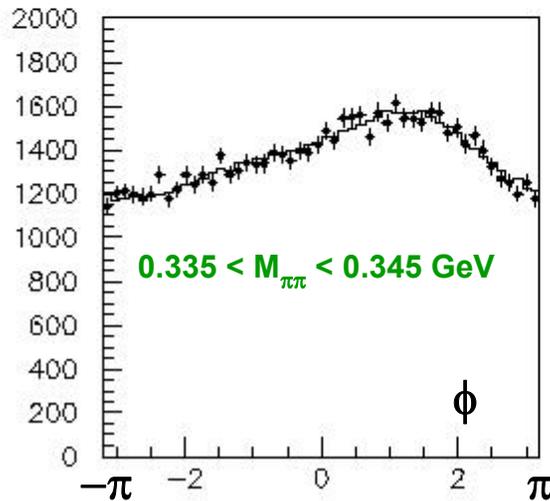
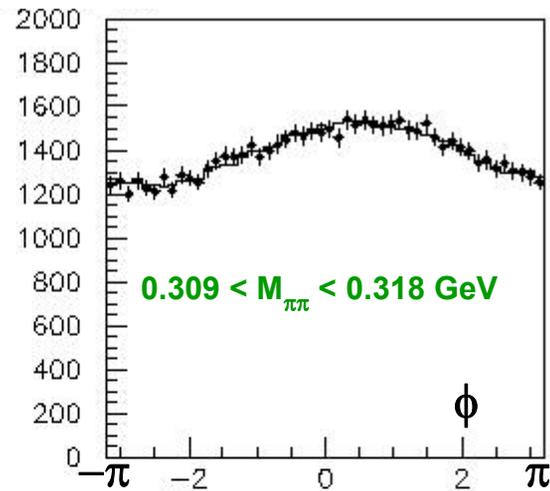
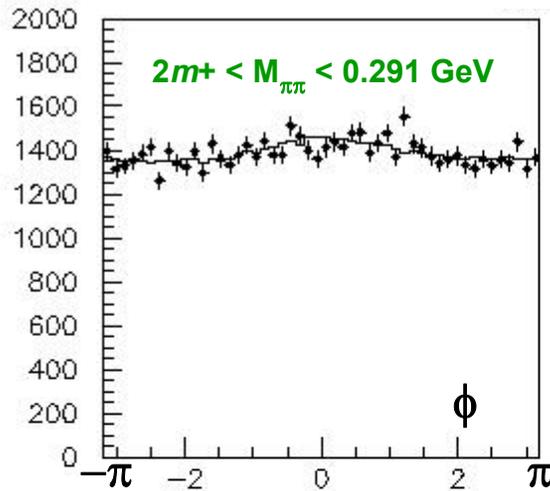
$\delta \neq 0$ in the Ke4 decay amplitude \longrightarrow asymmetric distribution of the charged lepton direction with respect to the plane defined by the two pions (ϕ distribution)

(Shabalin 1963)

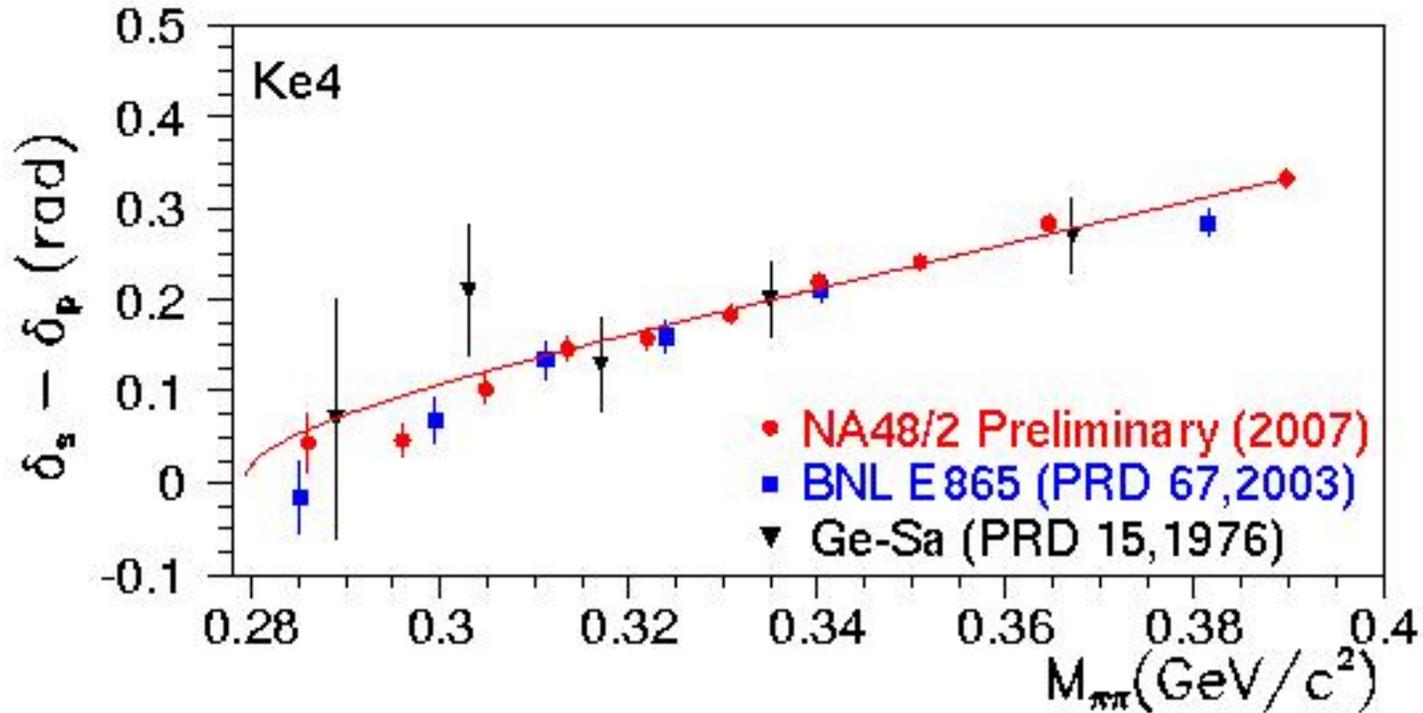
The asymmetry of the ϕ distribution increases with $M_{\pi\pi}$

NA48 / 2 (Preliminary)

ϕ distributions for four $M_{\pi\pi}$ bins



δ versus $M_{\pi\pi}$ (no isospin – breaking corrections)



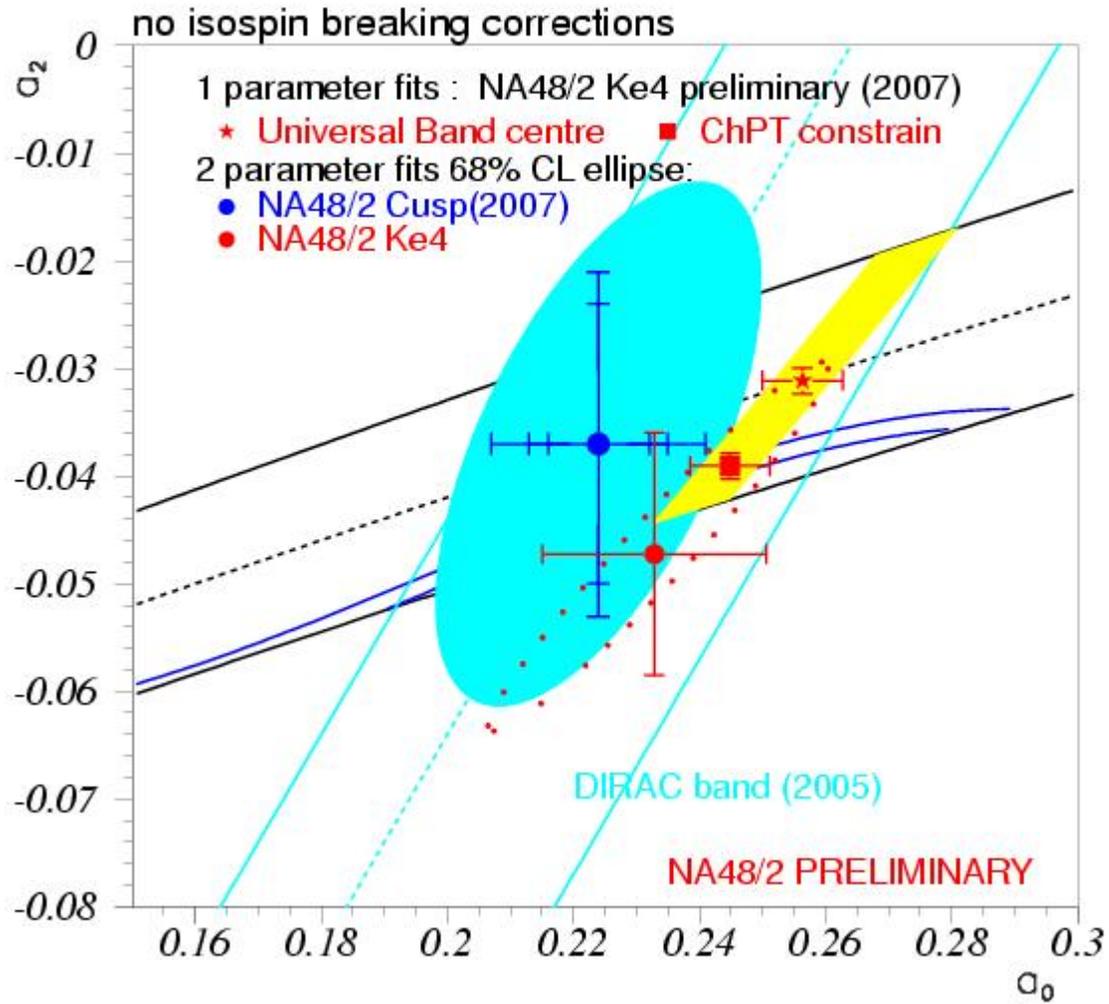
Geneva – Saclay : $\sim 30,000$ events , $p_{K^+} = 2.8 \text{ GeV}/c$

BNL E865 : 406,103 events (with $\sim 4.4\%$ background), $p_{K^+} = 6 \text{ GeV}/c$

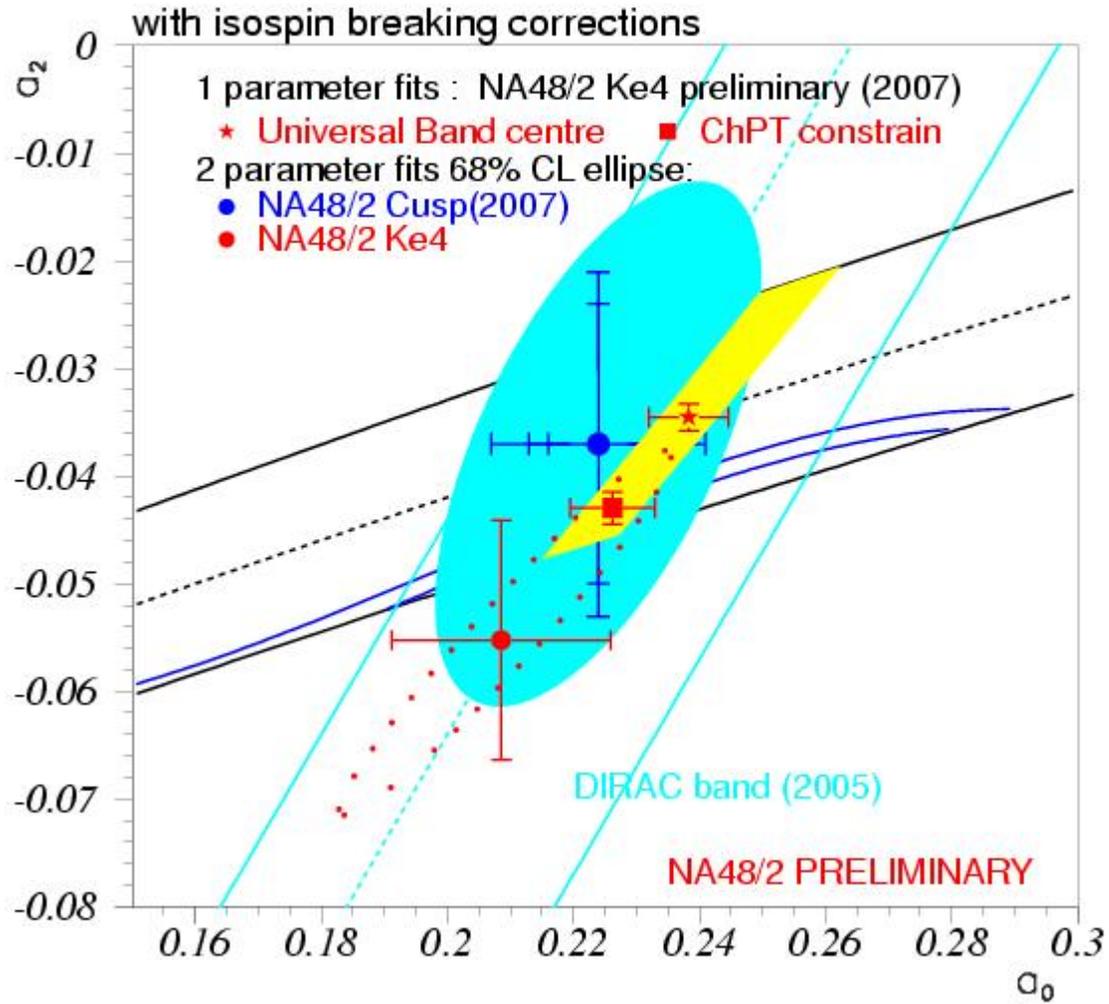
NA48/2 : 677,510 events (with $\sim 0.5\%$ background), $p_{K^\pm} = 6 \text{ GeV}/c$

NOTE: the isospin – breaking corrections reduce δ by 0.01 – 0.012
(see the presentation by J. Gasser)

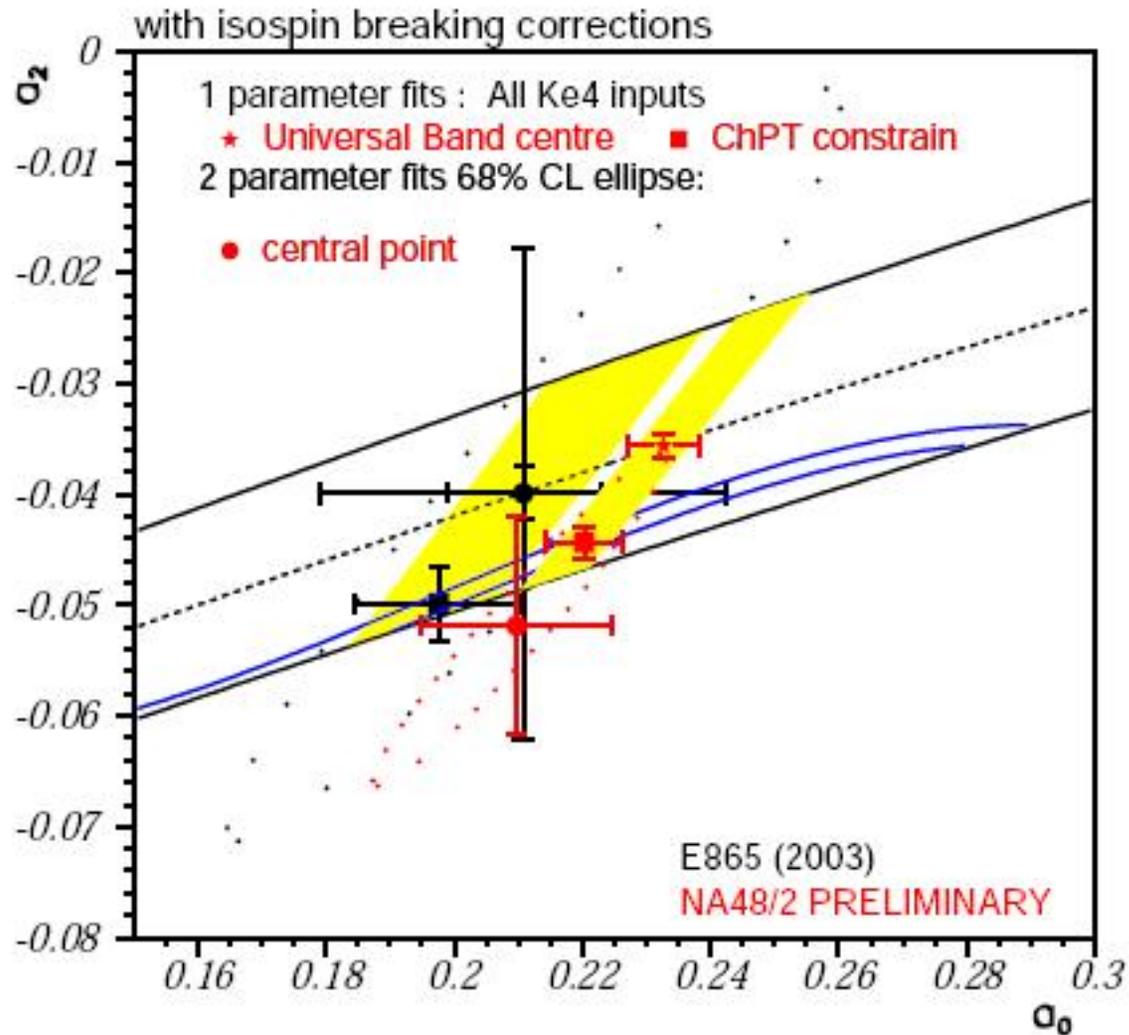
NA48/2 Cusp – Ke4 comparison



NA48/2 Cusp – Ke4 comparison



Ke4: NA48/2 – BNL E865 comparison



Marginal agreement for constrained fits

CONCLUSIONS

▪ Cusp in $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ decay

Little improvement is expected from the experimental side.

The determination of a_0 and a_2 can become more precise only if the theoretical models to fit the data are improved (electromagnetic effects, higher order terms)

Using the Bern – Bonn non-relativistic effective field theory to simultaneously fit the $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ and $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ data has just started.

A direct measurement of A_{+-} / A_{+00} from the NA48/2 data is very difficult, but (perhaps) not impossible. Will try.

▪ Cusp in $K_L \rightarrow \pi^0 \pi^0 \pi^0$ decay

Visible, but the effect is not large. Probably only a consistency check, not competitive with $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ to extract a precise value of $a_0 - a_2$.

▪ Ke4

Expect further improvement from NA48/2 by the addition of the 2004 data (\Rightarrow total number of events $> 10^6$).

A new development of the last few weeks: isospin – breaking corrections. Better agreement between NA48/2 Ke4 scattering lengths and those measured from the cusp.

Agreement between NA48/2 and E865 is marginal and should be understood